

## Studies of the Spin Hamiltonian Parameters and the Local Structure of Tetragonal $Zr^{3+}$ Centers in Orthophosphate $MPO_4$ ( $M = Sc, Lu, Y$ ) Crystals

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The spin Hamiltonian parameters ( $g$  factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$ ) of  $Zr^{3+}$  on the tetragonal  $M^{3+}$  ( $M = Sc, Lu, Y$ ) sites of zircon-type orthophosphate  $MPO_4$  crystals are calculated by high-order perturbation formulas of  $d^1$  ions in tetragonal symmetry. The crystal-field parameters are estimated by the superposition model and reasonable local structural data of impurity centers. The results show good agreement with the experimental values. It appears that in the case of size mismatch, the explanation of the spin Hamiltonian parameters of a paramagnetic impurity in crystals should take the impurity-induced local lattice relaxation into account.

**Key words:** EPR; Spin Hamiltonian Parameters; Crystal-Field Theory; Defect Structure;  $Zr^{3+}$ ;  $MPO_4$  ( $M = Sc, Lu, Y$ ).

### 1. Introduction

The orthophosphates  $MPO_4$  ( $M = Sc, Lu, Y$ ) form an isomorphous series of compounds with zircon-type structure. They and other zircon-type compounds doped with rare-earth and transition metal ions have gained considerable interest because of their application in lasers and luminescent materials [1–4]. Many optical and EPR spectral studies of doped  $MPO_4$  crystals were made [5–8]. For example, the EPR spectra of the  $4d^1$  configuration ion  $Zr^{3+}$  in  $MPO_4$  crystals were measured and their spin Hamiltonian parameters ( $g$  factors  $g_{\parallel}$ ,  $g_{\perp}$  and hyperfine structure constants  $A_{\parallel}$ ,  $A_{\perp}$ ) were given in [8]. The  $Zr^{3+}$  ion in  $MPO_4$  crystals replaces the  $M^{3+}$  ion and is in an eight-fold co-

ordinated site with tetragonal ( $D_{2d}$ ) symmetry. Until now no theoretical explanations (in particular those related to the structure of impurity centers) for these spin Hamiltonian parameters were given. Abraham et al. [8] thought that these spin Hamiltonian parameters were not accounted for by previously published second-order perturbation expressions. In the present paper, we calculate these spin Hamiltonian parameters by high-order perturbation formulas of the  $d^1$  ion in tetragonal symmetry. In the calculations, the crystal-field parameters are estimated by the superposition model [9] and reasonable local structure data of impurity centers. The results are discussed.

### 2. Calculation

For a  $d^1$  configuration ion, a cubic crystal-field will split the five-fold degeneracy energy level  $^2D$  into an orbital doublet  $^2E(D)$  and triplet  $^2T_2(D)$ . In the eight-fold coordinated site under study, the energy level  $^2E(D)$  is lower. Adding the tetragonal distortion, the  $^2E(D)$  level is split into two different orbital singlets  $^2A_1$  and  $^2B_1$ , and the  $^2T_2(D)$  is split into a doublet  $^2E(T)$  and a singlet  $^2B_2$ . From the observed  $g$  factors of  $MPO_4:Zr^{3+}$  [8] one finds that the ground state is  $^2B_1(|d_{x^2-y^2}\rangle)$ . By the spin Hamiltonian theory, the third-order perturbation formulas of  $g_{\parallel}$ ,  $g_{\perp}$ ,  $A_{\parallel}$  and  $A_{\perp}$  for a  $d^1$  ion in tetragonal symmetry can be expressed as

$$\begin{aligned} g_{\parallel} &= g_s - \frac{8k\zeta}{E_1} - \frac{(g_s + k)\zeta^2}{E_2^2} - \frac{4k\zeta^2}{E_1E_2}, \\ g_{\perp} &= g_s - \frac{2k\zeta}{E_2} - \frac{(g_s/2 - k)\zeta^2}{E_2^2} - \frac{2g_s\zeta^2}{E_1^2}, \\ A_{\parallel} &= P \left[ -K - \frac{4}{7} + (g_{\parallel} - g_s) + \frac{3}{7}(g_{\perp} - g_s) \right], \\ A_{\perp} &= P \left[ -K + \frac{2}{7} + \frac{11}{14}(g_{\perp} - g_s) \right], \end{aligned} \quad (1)$$

where  $g_s$  ( $\approx 2.0023$ ) is the free-electron value.  $\zeta$  and  $P$  are the spin-orbit coupling parameter and dipolar hyperfine structure constant in the crystal.  $k$  and  $K$  are the orbital reduction factor and the core polarization constant. Because of the covalence reduction effect for  $d^n$  ions in crystals, we have [10–12]

$$\zeta \approx N^2\zeta_0, \quad P \approx N^2P_0,$$

	$R_1^h$ (Å) [21]	$R_2^h$ (Å) [21]	$R_1$ (Å)	$R_2$ (Å)	$\bar{R}$ (Å)	$\theta_1$ (deg) [21]	$\theta_2$ (deg) [21]	$k$
ScPO <sub>4</sub> :Zr <sup>3+</sup>	2.150	2.277	2.229	2.356	2.293	103.17	31.62	0.911
LuPO <sub>4</sub> :Zr <sup>3+</sup>	2.264	2.346	2.284	2.366	2.325	103.47	30.95	0.934
YPO <sub>4</sub> :Zr <sup>3+</sup>	2.313	2.374	2.3115	2.3725	2.342	103.67	30.22	0.944

Table 1. The structural data and the orbital reduction factor  $k$  for MPO<sub>4</sub>:Zr<sup>3+</sup> (M = Sc, Lu, Y) crystals.

Table 2. The spin Hamiltonian parameters of MPO<sub>4</sub>:Zr<sup>3+</sup> (M = Sc, Lu, Y) crystals.

		ScPO <sub>4</sub> :Zr <sup>3+</sup>	LuPO <sub>4</sub> :Zr <sup>3+</sup>	YPO <sub>4</sub> :Zr <sup>3+</sup>
$g_{\parallel}$	Cal. <sup>a</sup>	1.892	1.849	1.830
	Cal. <sup>b</sup>	1.870	1.844	1.830
	Expt. [8]	1.871	1.844	1.832
$g_{\perp}$	Cal. <sup>a</sup>	1.948	1.936	1.933
	Cal. <sup>b</sup>	1.938	1.934	1.933
	Expt. [8]	1.936	1.933	1.932
$A_{\parallel}$ (10 <sup>-4</sup> cm <sup>-1</sup> )	Cal. <sup>a</sup>	91.5	96.3	98.3
	Cal. <sup>b</sup>	92.8	96.6	98.3
	Expt. [8]	92.7	98.5	100.8
$A_{\perp}$ (10 <sup>-4</sup> cm <sup>-1</sup> )	Cal. <sup>a</sup>	46.2	47.9	48.5
	Cal. <sup>b</sup>	46.6	47.9	48.5
	Expt. [8]	49.9	52.1	52.9

<sup>a</sup> Calculated using the structural data of host crystals.

<sup>b</sup> Calculated, using the defect structural data of impurity centers.

in which  $\zeta_0$  and  $P_0$  are the corresponding values of a free  $d^n$  ion, and  $N^2$  ( $\approx k$ ) is the covalence reduction factor. For a free Zr<sup>3+</sup> ion,  $\zeta_0 \approx 500$  cm<sup>-1</sup> [13].  $P_0 \approx g_e g_N \beta_e \beta_N \langle r^{-3} \rangle$  [14], and from the value of  $\langle r^{-3} \rangle \approx 3.16$  a. u. [13] we obtain  $P_0 \approx -52.5 \cdot 10^{-4}$  cm<sup>-1</sup> for a free <sup>91</sup>Zr<sup>3+</sup> ion. The energy separations  $E_1$  and  $E_2$  can be written as

$$\begin{aligned} E_1 &= E(^2B_2) - E(^2B_1) = 10D_q, \\ E_2 &= E(^2E_2) - E(^2B_1) = 10D_q - 3D_s + 5D_t. \end{aligned} \quad (2)$$

The crystal-field parameters  $D_q$ ,  $D_s$  and  $D_t$  of the studied system in the superposition model [9] can be given as

$$\begin{aligned} D_q &= \frac{4}{3} \bar{A}_4(R_0) \sum_{i=1}^2 \left[ \left( \frac{R_0}{R_i} \right)^{t_4} \sin^4 \theta_i \right], \\ D_s &= -\frac{4}{7} \bar{A}_2(R_0) \sum_{i=1}^2 \left[ \left( \frac{R_0}{R_i} \right)^{t_2} (3 \cos^2 \theta_i - 1) \right], \\ D_t &= -\frac{4}{21} \bar{A}_4(R_0) \sum_{i=1}^2 \left[ \left( \frac{R_0}{R_i} \right)^{t_4} (7 \sin^4 \theta_i + 35 \cos^4 \theta_i \right. \\ &\quad \left. - 30 \cos^2 \theta_i + 3) \right], \end{aligned} \quad (3)$$

where the power-law exponents are  $t_2 \approx 3$  and  $t_4 \approx 5$  because of the ionic nature of the bonds [9, 15, 16].  $\bar{A}_2(R_0)$  and  $\bar{A}_4(R_0)$  are the intrinsic parameters depending on the ligand, the coordination number and

reference distance  $R_0$ . The values of  $\bar{A}_4(R_0)$  and the ratio  $\bar{A}_2(R_0)/\bar{A}_4(R_0) \approx 8 \sim 12$  were obtained for many  $3d^n$  MX<sub>m</sub> clusters in crystals from optical and EPR spectra studies [15–18]. However, for the  $4d^n$  MX<sub>m</sub> clusters in crystals, few values of  $\bar{A}_4(R_0)$  and  $\bar{A}_2(R_0)/\bar{A}_4(R_0)$  were reported. From the optical spectra of similar [ZrO<sub>8</sub>]<sup>13-</sup> clusters in a YAG:Zr<sup>3+</sup> crystal [19] we estimate  $\bar{A}_4(R_0) \approx 1750$  cm<sup>-1</sup> with  $R_0 \approx 2.3$  Å. It can be expected that the ratio  $\bar{A}_2(R_0)/\bar{A}_4(R_0)$  for  $4d^n$  clusters is smaller than that for  $3d^n$  clusters because (i) the ratio  $\bar{A}_2(R_0)/\bar{A}_4(R_0)$  is related to the ratio  $\langle r_2 \rangle / \langle r_4 \rangle$  and (ii) the ratio  $\langle r_2 \rangle / \langle r_4 \rangle$  of a  $4d^n$  ion is smaller than that of a  $3d^n$  ion [13] (e. g.,  $\langle r_2 \rangle / \langle r_4 \rangle \approx 0.15$  and  $0.27$  a. u. for  $4d^1$  Zr<sup>3+</sup> and  $3d^1$  Ti<sup>3+</sup> ions, respectively [13]). So, we take  $\bar{A}_2(R_0)/\bar{A}_4(R_0) \approx 5.4$  for MPO<sub>4</sub>:Zr<sup>3+</sup> crystals.  $R_i$  ( $i = 1$  or  $2$ ) is the impurity-ligand distance and  $\theta_i$  is the angle between the  $R_i$  and  $C_4$  axis. In general, the local structural data of impurity centers are unlike the corresponding data in the host crystal because of the size and/or charge mismatch. For the studied MPO<sub>4</sub>:Zr<sup>3+</sup>, we can reasonably let the angle  $\theta_i$  be unchanged and estimate the impurity-ligand distance  $R_i$  from the approximate formula [20]

$$R_i \approx R_i^h + (r_i - r_h)/2, \quad (4)$$

where  $R_i^h$  is the corresponding distance in the host crystal. The values of  $R_i^h$  for MPO<sub>4</sub> crystals [21] are shown in Table 1.  $r_i$  and  $r_h$  are the ionic radius of the impurity and that of the replaced host ion. In MPO<sub>4</sub>:Zr<sup>3+</sup> crystals,  $r_h \approx 0.732$ ,  $0.85$  and  $0.892$  Å [22] for M<sup>3+</sup> = Sc<sup>3+</sup>, Lu<sup>3+</sup> and Y<sup>3+</sup>, respectively, and  $r_i$  (Zr<sup>3+</sup>)  $\approx 0.89$  Å, estimated by the interpolation method based on  $r(\text{Zr}^+) \approx 1.09$  Å and  $r(\text{Zr}^{4+}) \approx 0.79$  Å [22]. Thus, for MPO<sub>4</sub>:Zr<sup>3+</sup> crystals the impurity-ligand distances  $R_i$  are estimated. They and the angles  $\theta_i$  are collected in Table 1. So, in the above formulas only the parameters  $K$  and  $k$  are not known. By fitting the calculated  $g_{\parallel}$ ,  $g_{\perp}$ ,  $A_{\parallel}$  and  $A_{\perp}$  to the observed values, we obtain the same value  $K \approx 1.21$  and different  $k$  values (see Table 1) for the three MPO<sub>4</sub>:Zr<sup>3+</sup> crystals. The calculated spin Hamiltonian parameters are compared with the experimental values in Table 2.

In addition, for comparison, we also calculate the spin Hamiltonian parameters for  $\text{MPO}_4:\text{Zr}^{3+}$  crystals by using the structural data in the host crystals. The calculated results are also shown in Table 2.

### 3. Discussion

It is agreed that for a  $\text{MX}_m$  cluster the covalence increases and hence the covalence reduction factor  $N^2(\approx k)$  decreases with the decreasing M-X distance  $R$  [10, 23, 24]. Considering the different average  $\text{Zr}^{3+}-\text{O}^{2-}$  distances  $\bar{R}$  of  $[\text{ZrO}_8]^{13-}$  clusters in  $\text{MPO}_4:\text{Zr}^{3+}$  crystals (see Table 1), one should expect an increasing value of  $k$  with an increase of the distance  $\bar{R}$ ,  $k(\text{ScPO}_4:\text{Zr}^{3+}) < k(\text{LuPO}_4:\text{Zr}^{3+}) < k(\text{YPO}_4:\text{Zr}^{3+})$ . Our calculated values of  $k$  (see Table 1) are consistent with this expectation and can be regarded as suitable.

The calculated spin Hamiltonian parameters for  $\text{MPO}_4:\text{Zr}^{3+}$  crystals, using the defect structural data of impurity centers, show good agreement with the experimental values (see Table 2). If the structural data of the host crystals are applied, the calculated results of  $\text{LuPO}_4:\text{Zr}^{3+}$  and  $\text{YPO}_4:\text{Zr}^{3+}$  are close to the observed values, but the calculated  $g_{\parallel}$  and  $g_{\perp}$  for  $\text{ScPO}_4:\text{Zr}^{3+}$  agree poorly with the observed values, because the ionic radius of the impurity  $\text{Zr}^{3+}$  is close to those of the replaced host ions in  $\text{LuPO}_4:\text{Zr}^{3+}$  and  $\text{YPO}_4:\text{Zr}^{3+}$ , but the difference of the ionic radius between  $\text{Zr}^{3+}$  and the replaced  $\text{Sc}^{3+}$  ion in the  $\text{ScPO}_4:\text{Zr}^{3+}$  crystal is larger. So, in the cases of size or/and charge mismatch, the reasonable explanations of spin Hamiltonian parameters for a paramagnetic impurity in crystals should take the impurity-induced local lattice relaxation into account.

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