

The D_{2h} Distortion around the Cu^{2+} Center in $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ Single Crystals

Y. Huang, M. L. Du, C. Ni, and J. F. Wen

Department of Physics, Southwest University for Nationalities, Chengdu 610041, P. R. China

Reprint requests to Dr. M. L. D.; E-mail: duml@mail.sc.cninfo.net

Z. Naturforsch. **60a**, 193 – 195 (2005); received September 16, 2004

A formula for the calculation of the three g factors of $3d^9$ ions in an orthorhombic field D_{2h} has been derived. Using it to investigate the EPR g factors of the Cu^{2+} ions in single crystals of $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$, the variation of the g factors on changing the angle α between the x - and y -axis has been explained. According to that, it can be confirmed that the angle α of the D_{2h} distortion is about 62.6° . PACS: 71.70C; 76.30F

Key words: $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ Crystal; Gyromagnetic Factor; D_{2h} Distortion.

1. Introduction

$\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystals belong to the Nasicon-type family with a three-dimensional network built of PO_4 tetrahedra sharing corners with ZrO_6 octahedra [1]. The three-dimensional network can be considered as being made of infinite ribbons linked by PO_4 tetrahedra. They are used in chemistry and ceramic industry because of their catalytic and low thermal expansion properties, as well the ionic conductivity of their derivatives [2–8].

Taoufik et al. [9] have investigated the magnetic susceptibility and EPR of $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystals. They contain an important amount of paramagnetic Cu^{2+} ions, and EPR spectra give information about local paramagnetic environments. Their structure shows a monoclinic distortion compared to that of $\text{NaZr}_2(\text{PO}_4)_3$ [6]. It is suggested that the field surrounding the Cu^{2+} ions is orthorhombic (D_{2h}) rather than tetragonal (D_{4h}) [9, 10] from the observed optical spectrum of the Cu^{2+} ions. But these observations didn't confirm the distortion structure when the crystal field varies from D_{4h} to D_{2h} .

In this paper, using experimental EPR results, further studies have been done to observe the distortion tendency from D_{4h} to D_{2h} .

2. The g Factors of $3d^9$ in the Symmetry of D_{2h}

The Cu^{2+} ions lie in the interspace of the three-dimensional network and are surrounded by six oxygen

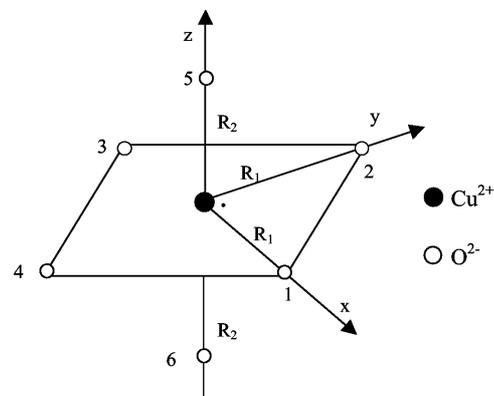


Fig. 1. Position of the six ligands in D_{2h} .

atoms [2]. Taoufik et al. assumed that it is the D_{2h} symmetry [9] shown in Fig. 1, where the angle α between the x - and y -axis lies in the plane perpendicular to the z -axis. For the $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystal $R_1 = 1.95 \text{ \AA}$ and $R_2 = 2.82 \text{ \AA}$ have been given in [2], but the resulting angle α was not confirmed.

It is known that Cu^{2+} belongs to the electron system $3d^9$. Its energy level will be split into 2E and 2T_2 in a cubic field. The ground state is 2E in octahedral symmetry. In the orthorhombic field D_{2h} the energy levels will be split further. 2T_2 is split into $B_1(\zeta)$, $B_2(\eta)$, and $B_3(\xi)$. 2E is split into $A_1(\varepsilon)$ and $A_1(\theta)$. A_1 , B_1 , B_2 , and B_3 are the irreducible representation in D_{2h} symmetry. ε and θ indicate the components of 2E . ζ , η , and ξ indicate the components of 2T_2 . $A_1(\varepsilon)$ is the ground state in the $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystal [11].

In an orthorhombic field, the spin Hamiltonian of the $3d^9$ ion can be described by the expression

$$H_s = g_x \mu_B H_x \hat{S}_x + g_y \mu_B H_y \hat{S}_y + g_z \mu_B H_z \hat{S}_z, \quad (1)$$

where g_i ($i = x, y, z$) indicates the components of the g factor, μ_B is the Bohr magneton, \hat{S}_i ($i = x, y, z$) is the spin operator, and H_i ($i = x, y, z$) indicates the components of the magnetic field along the x -, y - and z -axes.

Using the perturbation theory, the g factors can be obtained by the formula [12]

$$g_i = g_s - 2\lambda A_{ii}, \quad (2)$$

$$A_{ij} = \sum_{n \neq 0} \frac{\langle 0 | \hat{L}_i | n \rangle \langle n | \hat{L}_j | 0 \rangle}{E_n^{(0)} - E_0^{(0)}}, \quad (i, j = x, y, z), \quad (3)$$

where $g_s = 2.0023$ is the value of free electron and λ is the spin-orbit coupling coefficient of the $3d^9$ ion. The relation between λ and the one-electron spin-orbit coupling coefficient ζ_d is $\lambda = -\zeta_d$. Using (2) and introducing the average covalent factor N [13] to describe the covalency, the g factors of the $3d^9$ ion in D_{2h} symmetry can be obtained as

$$g_z = g_s - \frac{8\zeta_d N^4}{E(\zeta) - E(\varepsilon)}, \quad (4)$$

$$g_x = g_s - \frac{2\zeta_d N^4}{E(\xi) - E(\varepsilon)}, \quad (5)$$

$$g_y = g_s - \frac{2\zeta_d N^4}{E(\eta) - E(\varepsilon)}, \quad (6)$$

where the energy denominators are

$$E(\zeta) - E(\varepsilon) = -\frac{2}{3} \sqrt{\frac{10}{7}} B_{44}, \quad (7)$$

$$E(\eta) - E(\varepsilon) = -\frac{1}{3} \sqrt{\frac{10}{7}} B_{44} + \frac{2\sqrt{10}}{21} B_{42} - \frac{2}{21} B_{40} + \frac{\sqrt{6}}{7} B_{22} + \frac{2}{7} B_{20}, \quad (8)$$

$$E(\xi) - E(\varepsilon) = -\frac{1}{3} \sqrt{\frac{10}{7}} B_{44} - \frac{2\sqrt{10}}{21} B_{42} - \frac{2}{21} B_{40} - \frac{\sqrt{6}}{7} B_{22} + \frac{2}{7} B_{20}. \quad (9)$$

The crystal-field parameter B_{kq} is related to the crystal structure parameter. In D_{2h} symmetry the crystal-field parameters B_{44} , B_{42} , B_{40} , B_{22} , and B_{20} are related

to the band lengths R_1 , R_2 , and the angle α . They can be obtained from the expressions

$$B_{44} = B_{4-4} = \frac{-1}{4} \sqrt{\frac{25}{2}} (1 + \cos 4\alpha) \frac{eq}{R_1^5} \langle r^4 \rangle, \quad (10)$$

$$B_{42} = B_{4-2} = \frac{1}{2} \sqrt{\frac{5}{2}} (1 + \cos 2\alpha) \frac{eq}{R_1^5} \langle r^4 \rangle, \quad (11)$$

$$B_{40} = \frac{-1}{2} \left(\frac{3}{R_1^5} + \frac{4}{R_2^5} \right) eq \langle r^4 \rangle, \quad (12)$$

$$B_{22} = B_{2-2} = -\sqrt{\frac{3}{2}} (1 + \cos 2\alpha) \frac{eq}{R_1^3} \langle r^2 \rangle, \quad (13)$$

$$B_{20} = -2 \left(\frac{1}{R_2^3} + \frac{1}{R_2^3} \right) eq \langle r^2 \rangle, \quad (14)$$

where q is the charge of the ligand, e the charge of the electron, and $\langle r^2 \rangle$ and $\langle r^4 \rangle$ are the expectation values in the crystal.

3. The D_{2h} Distortion Structure

Considering the average covalent factor N , the relations between the expectation values $\langle r^k \rangle$ in the crystal and $\langle r^k \rangle_0$ in the free ion are

$$\langle r^2 \rangle = N^2 \langle r^2 \rangle_0, \quad \langle r^4 \rangle = N^2 \langle r^4 \rangle_0, \quad (15)$$

and the relation between the spin-orbit coupling coefficient ζ_d in the crystal and ζ_d^0 in the free ion is

$$\zeta_d = N^2 \zeta_d^0. \quad (16)$$

The expectation values $\langle r^k \rangle_0$ in a free Cu^{2+} ion are [14]

$$\langle r^2 \rangle_0 = 3.11 a_0^2, \quad \langle r^4 \rangle_0 = 44.80 a_0^4, \quad (17)$$

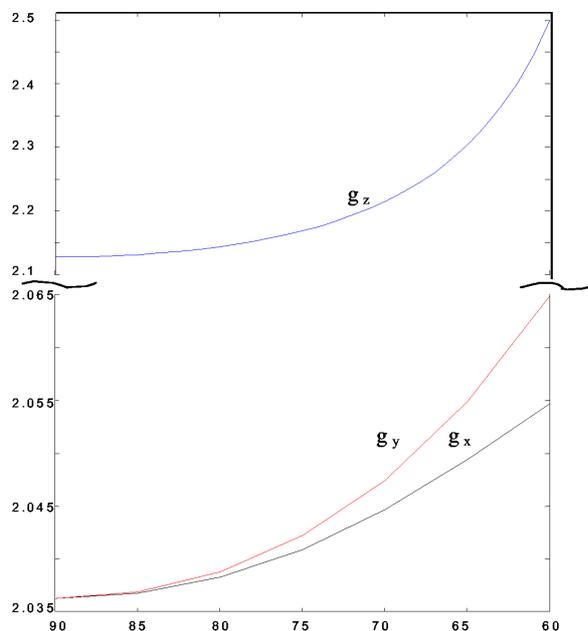
where a_0 is the Bohr radius. The spin-orbit coupling coefficient ζ_d^0 in the free Cu^{2+} is [15]

$$\zeta_d^0 = 829 \text{ cm}^{-1}. \quad (18)$$

Then the g factors can be calculated. From (10)–(14), the g factors depend on the band lengths R_1 , R_2 , and the angle α . When α is 90° , g_x and g_y are equal. When α is not 90° , the difference between g_x and g_y is not zero. Moreover, we can get the varying tendency of the g factors with the angle α , when the crystal field around Cu^{2+} is distorted from D_{4h} to D_{2h} . The results are shown in Figure 2.

Table 1. Comparison of the theoretical and experimental results ($\alpha = 62.6^\circ$, $N = 0.98$)

g Factor	Calculated	Experimental
g_x	2.0519	2.068
g_y	2.0594	2.071
g_z	2.3772	2.374

Fig. 2. Dependence of the components of the g factor on the angle α ($N = 0.98$).

As shown in Fig. 2, when the angle α decreases from 90° , g_z increases. At $\alpha = 90^\circ$ i. e. when the field is D_{4h} , g_x and g_y are equal. When α decreases from 90° ,

g_x and g_y increase, whereby g_y increases more than g_x . According to the EPR experimental data, g_x and g_y are different. It shows that the crystal field around the center of the Cu^{2+} ion is D_{2h} . This confirms Taoufik's analysis [9]. From (4)–(6), the values of g_x , g_y , and g_z are related to the crystal structure data R_1 , R_2 , and α . Taking the angle α as the fitting parameter, we can fit the experimental values of the g factors ($g_z = 2.374$, $g_x = 2.068$, $g_y = 2.071$) [9]. The results are shown in Table 1. The theoretical values are very close to the experimental ones, when α is about 62.6° .

Thereby it is reasonable and satisfactory to explain the paramagnetic g factors of Cu^{2+} ions in $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystals. The crystal field around the central Cu^{2+} ion is D_{2h} indeed. From the EPR experiment, the angle α is about 62.6° .

4. Conclusion

In this paper, formulas for the calculation of the three g factors of $3d^9$ ions in an orthorhombic field D_{2h} have been given. They are related to the angle α and the band lengths R_1 and R_2 of the crystal structure. With these formulas, the tendency of the g factors to vary with the angle α has been explained for $\text{Cu}_{0.5}\text{Zr}_2(\text{PO}_4)_3$ crystal. The best fitting value of the angle α is 62.6° .

Acknowledgement

The project is supported by the National Science Foundation of China (Grant No. 50372053).

- [1] L. O. Hagman and P. Kierkegaard, *Acta Chem. Scand.* **22**, 1822 (1968).
- [2] I. Bussereau, R. Olazcuaga, G. Le Flem, and P. Hagenmuller, *Eur. J. Solid State Inorg. Chem.* **26**, 383 (1989).
- [3] E. Fargin, I. Bussereau, R. Olazuaga, G. Le Flem, C. Cartier, and H. Dexpert, *J. Solid State Chem.* **112**, 176 (1994).
- [4] A. El Jazouli, *Adv. Mater. Res.* **1–2**, 105 (1994).
- [5] I. Bennouna, S. Arsalane, R. Brochu, M.R. Lee, J. Chassaing, and M. Quarton, *J. Solid State Chem.* **112**, 224 (1994).
- [6] A. Serghini, R. Brochu, M. Ziyad, and J.C. Vedrine, *J. Alloys Compd.* **188**, 60 (1992).
- [7] A. El Jazouli, M. Alami, R. Brochu, J. M. Dance, G. Le Flem, and O. Hagenmuller, *J. Solid State Chem.* **71**, 444 (1987).
- [8] A. Serghini, R. Brochu, M. Ziyad, M. Loukah, and J. C. Vedrine, *J. Chem. Soc. Faraday Trans.* **87** (15), 2487 (1991).
- [9] I. Taoufik, M. Haddad, A. Nadiri, R. Brochu, and R. Berger, *J. Phys. Chem. Solids.* **60**, 701 (1999).
- [10] R. Debnath and J. Chaudhuri, *J. Phys. Chem. Glasses* **36** (4), (1995) 160.
- [11] I. S. Ahuja and S. Tripathi, *Spectrochim. Acta Part A:* **47** (5), (1991) 637.
- [12] M. G. Zhao, *Ligand Field Theory*, Guizhou People Press, Guiyang (China), 328 (1985) (in Chinese).
- [13] M. G. Zhao, J. A. Xu, G. R. Bai, and H. S. Xie, *Phys. Rev. B* **27**, 1516 (1983).
- [14] M. G. Zhao, Z. X. Hu, and Q. Z. Huang, *J. Geochem (China)* **1**, 44 (1979) (in Chinese).
- [15] I. H. Parker, *J. Phys. C* **4**, 2967 (1971).