

Anisotropic Anomalous Dispersion in Cuprite, Cu_2O

Klaus Eichhorn

Hamburger Synchrotron-Strahlungs-Labor HASYLAB

Armin H. Kirfel and Karl F. Fischer

Fachrichtung Kristallographie, Fachbereich Physik, Universität des Saarlandes Saarbrücken

Z. Naturforsch. **43a**, 391–392 (1988);

received December 23, 1987

The paper reports the observation of significant "resonance peaks" in the intensity vs. energy curves around the Cu-K-absorption edge for forbidden reflections in centrosymmetric cubic Cu_2O . A qualitative interpretation is given in terms of an anisotropic anomalous dispersion model based on f' , f'' -tensors compatible with the symmetry at the Cu-site.

Keywords: Synchrotron radiation; X-ray diffraction; Anomalous dispersion; X-ray absorption edge; Cuprite.

Violations of the "special" extinction rules caused by anisotropic thermal vibrations and/or by bond-induced deformations of the atomic valence shells are well known [1]. In addition, anisotropic behaviour of the anomalous dispersion around K- and L-absorption edges can give rise to observation of "forbidden" reflections, even to those which are invariant against anisotropies of the thermally smeared atomic scattering factor. The availability of synchrotron light with its tunability over a wide energy range enables detailed investigations of this effect.

In conventional X-ray work anomalous dispersion is taken into account by the energy-dependent correction terms f' and f'' [2]. Recent experiments and theoretical studies have shown that f' , f'' can also vary in reciprocal space due to a local anisotropy of the X-ray susceptibility in the crystal [3,4]. Consequently, in a limited energy range around an absorption edge the anomalous dispersion is not only an atomic property but it is also affected by the chemical and structural environment of the anomalous scatterer under consideration. Then, the next approximations for describing f' and f'' are second rank tensors, in general $f'(\mathbf{h}, E)$ and $f''(\mathbf{h}, E)$ at \mathbf{h} in reciprocal space, and their analysis is *no more* confined to acentric structures.

Cuprite, Cu_2O , crystallizes in the cubic space group $\text{Pn}\bar{3}\text{m}$, $a = 4.2696 \text{ \AA}$. Cu, in $2b(\bar{3}m, 0, 0, 0)$, is linearly bonded to two oxygen atoms in $4a(\bar{4}3m, 1/4, 1/4, 1/4)$. This arrangement provides a highly anisotropic Cu-environment, which is reflected in one of the largest known electric field gradients, at the Cu position $|\nabla E_{zz}| = 134 \cdot 10^{20} \text{ Vm}^{-2}$ [5,6].

Reflections of parity (eoo) are strictly forbidden for spherically symmetric atoms. However, anisotropic vibration of Cu

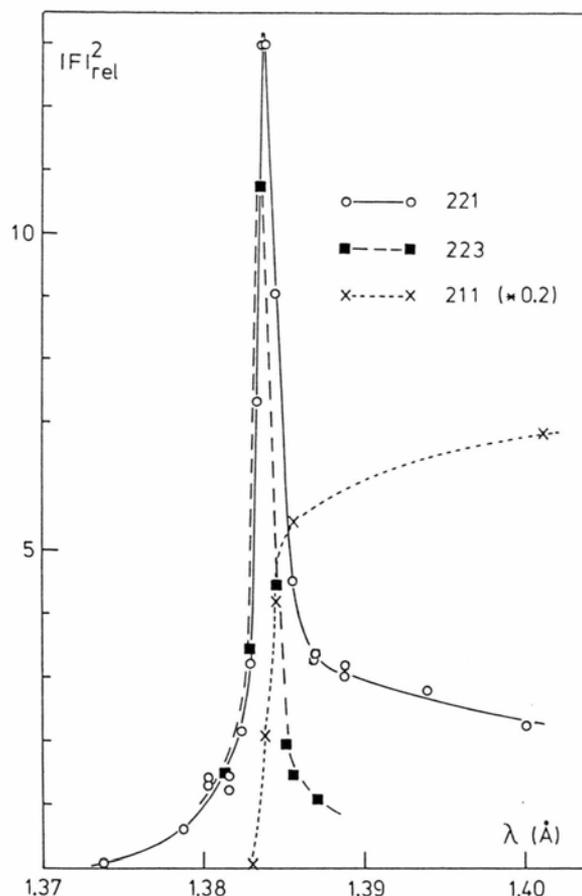


Fig. 1. Observed intensities of reflections (221), (223), and (211) as function of photon energy around the Cu-K-absorption edge (intensity factor for (211) is 0.02 instead of 0.2).

($U_{12} \neq 0$) and a quadrupolar deformation of its valence shell [5] give rise to small intensities for (eoo)-reflections. Any anisotropy of f' , f'' around the Cu-K-absorption edge will result in an additional contribution to $I(\text{eoo})$ and should be readily observable with an intense tunable X-ray source.

Forbidden reflections from a natural Cuprite crystal ($5.2 \times 5.4 \times 6.6 \text{ mm}^3$) have been investigated at room temperature on the Two-Axis-Diffractometer at HASYLAB/DESY [7]. The scattering plane was horizontal and parallel to the polarization plane of the incident beam. The photon energy was calibrated by an absorption spectrum taken on a $0.75 \mu\text{m}$ Cu foil. The intensities of two "forbidden" reflections, (221) and (223), and of the "allowed" (211) reflection were monitored at a series of photon energies around the Cu-K-absorption edge. (211) is a so called "oxygen" reflection which carries no information about the spherical part of the Cu electron density distribution. The intensities of the "forbidden" (221) and (223), Fig. 1, show impressive "resonance peaks" as function of energy, which *cannot* be ex-

Reprint requests to K. Eichhorn or A. Kirfel, DESY-HASYLAB, Notkestr. 85, D-2000 Hamburg 52, FRG.

0932-0784 / 88 / 0400-0391 \$ 01.30/0. – Please order a reprint rather than making your own copy.



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland Lizenz.

Zum 01.01.2015 ist eine Anpassung der Lizenzbedingungen (Entfall der Creative Commons Lizenzbedingung „Keine Bearbeitung“) beabsichtigt, um eine Nachnutzung auch im Rahmen zukünftiger wissenschaftlicher Nutzungsformen zu ermöglichen.

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

On 01.01.2015 it is planned to change the License Conditions (the removal of the Creative Commons License condition "no derivative works"). This is to allow reuse in the area of future scientific usage.

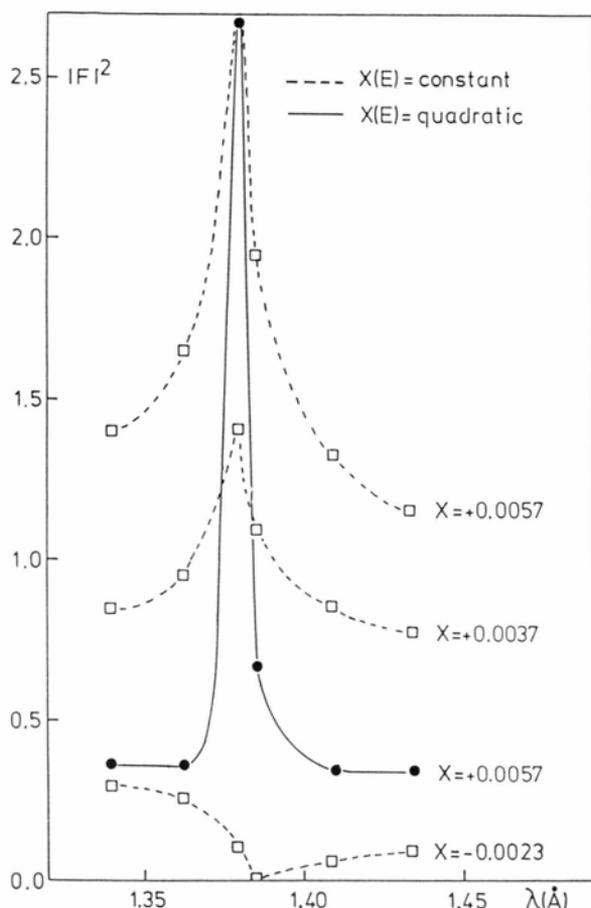


Fig. 2. Calculated intensities for (221) as function of photon energy. Solid line $X(E) = X_0[1 - |E - E_c|^{1/2}]$, broken lines $X(E) = \text{const}$.

- [1] B. Borie, *Acta Cryst.* **A37**, 238 (1981).
 [2] D. T. Cromer and D. A. Liberman, *J. Chem. Phys.* **53**, 1891 (1970).
 [3] D. H. Templeton and L. K. Templeton, *Acta Cryst.* **A38**, 62 (1982), *ibid* **A43**, 573 (1987).
 [4] V. E. Dmitrienko, *Acta Cryst.* **A39**, 29 (1983).

plained in terms of scalar f' , f'' correction factors. No significant changes were observed for the intensity of the "allowed" (211) apart from the expected absorption effect.

In order to explain these observations qualitatively we have modelled anisotropic f' , f'' by introducing an artificial anisotropic temperature factor for the K-shell with components $U_{11}(\text{Cu})$ and $U_{12}(\text{Cu}) + X$. Values $f'(\mathbf{h} = 0)$ and $f''(\mathbf{h} = 0)$ were taken from the measured Cu absorption spectra. This convolution with a Gaussian combined with an energy dependence, $X = X(E)$, of the anisotropy term gave a qualitative agreement between calculated and observed intensities $I_{221}(E)$ of the "resonance peak" (Fig. 2). This figure shows that X must be energy-dependent since the calculated intensity curves for constant X -values follow essentially the comparatively wide $f'(E)$ dispersion curve. In addition, there is a clear correlation of the calculated intensity curves with the sign of X . To fit the observations, X must be positive, rendering an anomalous dispersion "ellipsoid" elongated along [111]. This shape corresponds with that of the deformation potential distribution around Cu, whereas the thermal vibration of Cu is largest normal to the threefold axis ($U_{12} < 0$).

We conclude from the present results that we have observed directly the effect of anisotropic anomalous scattering in a centrosymmetric cubic structure as predicted by Dmitrienko [8]. Of course, there are anisotropic anomalous dispersion contributions to the intensities of all reflections. However, the effect will be difficult to detect in "normal" reflections due to the large contribution from the total electron density. Therefore, Forbidden Reflection near Edge Diffraction ("FRED") is of special interest, and more detailed studies of such reflections can yield information about the resonance behaviour and about the way the crystal field determines the anisotropy of f' , f'' . Resonance scattering may provide a local probe to the environment of the anomalous scatterer, and one may infer that this applies also to adsorbates on surfaces and to other two-dimensionally ordered systems.

The authors like to thank A. Petcov for help with the experiments. This work has received support by the Federal Minister of Research and Technology which is gratefully acknowledged.

- [5] R. Restori and D. Schwarzenbach, *Acta Cryst.* **B42**, 201 (1986).
 [6] H. Krüger and U. Meyer-Berkhout, *Z. Phys.* **132**, 171 (1952).
 [7] U. Bonse and K. Fischer, *Nucl. Instrum.* **190**, 4593 (1981).
 [8] V. E. Dmitrienko, *Acta Cryst.* **A40**, 89 (1984).